

Electronic properties of Pb-I deficient lead halide perovskites

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The electronic structure evolution of deficient halide perovskites with a general formula $(A, A')_{1+x}M_{1-x}X_{3-x}$ was investigated using the density functional theory. The focus is placed on characterization of changes in the band gap, band alignment, effective mass, and optical properties of deficient perovskites at various concentrations of defects. We uncover unusual electronic properties of the defect corresponding to a $M-X$ vacancy filled with an A' cation. This defect "repels" electrons and holes producing no trap states and, in moderate quantities ($x \leq 0.1$), does not hinder charge transport properties of the material. This behavior is rationalized using a confinement model and provides an additional insight to the defect tolerance of halide perovskites.

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I. INTRODUCTION

Recently, a new family of hybrid-halide perovskites called "deficient" perovskites (d-perovskites) and "hollow" perovskites was reported¹⁻⁶. Their structure is different from conventional AMX_3 perovskites where a relatively small organic cation A^+ , such as methyl ammonium ($A^+ = MA^+$), is centered within an inorganic MX_3^- cage formed by lead or tin halides. Deficient perovskites have missing $(M-X)^+$ units with corresponding voids filled by larger cations, such as hydroxyethylammonium ($A'^+ = HEA^+$), ethylenediammonium ($A'^+ = EDA^+$), or thioethylammonium ($A'^+ = TEA^+$). The defect (d) can be expressed using Kröger-Vink notations as

$$d^\times = v''_{Pb} + v^\bullet_I + (A')^\bullet_{Pb-I}. \quad (1)$$

The details of the defect reaction in Eq. (1) indicates that d consists of two intrinsic defects and one extrinsic defect. The superscript \times indicates null charge for the current site. v''_{Pb} is the first intrinsic defect which signifies the vacancy of Pb. The double prime $''$ denotes a negative charge (-2) due to missing of a Pb^{2+} . The second intrinsic defect is the vacancy of I v^\bullet_I , where the superscript \bullet denotes a net single positive charge ($+1$) due to missing of I^- . $(A')^\bullet_{Pb-I}$ is the extrinsic defect which demonstrates that A'^+ cation occupies the vacancies of Pb and I. This defect contributes a net single positive charge due to the presence of A'^+ .

The general formula of d-perovskites is $(A, A')_{1+x}M_{1-x}X_{3-x}$ ¹. The concentration x of such structural defects can be quite large (the range of $0 < x < 0.2$ has been demonstrated experimentally¹) leading to a range of $3 < X/M < 3.5$ that captures a deviation from conventional stoichiometry of perovskites. The difference between deficient and hollow perovskites comes from an ordered vs stochastic arrangement of defects, respectively. Despite a large concentration of structural defects, the material retains its 3D architecture and can be considered as a transition between the 2D and 3D structures⁷.

The most noticeable change in electronic properties of d-perovskites is the opening of an optical band gap. The band gap of d-MAPbI₃ with I/Pb ratio of 3.5 ($x = 0.2$) increases from 1.55 to 2.15 eV¹. Spanopoulos *et al.*³ also reported increase of the band gap from 1.52 to 1.85 eV for hollow MAPbI₃ perovskites (I/Pb ratio = 3.6) and from 1.25 to 1.42 eV for hollow MASnI₃ perovskites (I/Sn ratio = 3.6). Discrepancies in the experimental data can be attributed to different large cations (HEA^+ vs EDA^+). One should also note that the d-perovskites have been reported to present a dual band gap marked by the existence of

two shoulders in the UV/vis spectra [see Fig. 2(a) in Ref. 1] that leads to another uncertainty. Other noticeable properties of deficient/hollow perovskites are improved stability^{2,3,6}, improved photoluminescence efficiency of hollow MASnI_3 ⁴, and improved performance of Sn-based perovskite solar cells⁵.

If the issue of defects in halide perovskites and their impact on photovoltaic or light emission performances has been intensively investigated (see for instance Refs. 8,9 and references herein), theoretical studies of d-perovskites are scarce and limited to Refs. 3,5,6. *Ab initio* electronic structure calculations corroborate experimentally observed opening of the band gap; the band structure remains direct even at a high concentration of defects. It was also noticed that bands become less dispersed in hollow and d-perovskites^{3,6}. The valence band dispersion in a lateral direction (i.e., perpendicular to defect tunnels) practically vanished in d-FAPbI₃ [Fig. 5(b) in Ref. 6] at a high concentration of defects ($x = 0.2$) that corresponds to formation of a continuous chain of defects (tunnels). Leblanc *et al.*⁶ also pointed that the band gap opening with increasing x occurs mostly due to a shift of the conduction band edge (CBE) upward in energy, while the valence band edge (VBE) shifts down by approximately 0.1 eV. Similar to semiconductor heterojunctions induced band offset, defects will introduce energy misalignment. Fig. 1 illustrated two different types of band offsets, namely the type I band alignment or straddling gap and the type II band alignment or staggered gap. Identification of the type of band alignment is critical for investigating carrier transport properties in semiconductors.

Here, we use first-principle calculations to investigate the electronic structure evolution of d-MAPbI₃. Two deficient structures are selected with the ratio I/Pb = 3.22 and 3.5 to represent moderate ($x = 0.1$) and high density of defects ($x = 0.2$), respectively. We quantify changes in the band gap, band alignment, effective mass, and effective band structure of d-MAPbI₃ based on the structural model from Ref. 1. We anticipate the results to provide further insight into understanding a defect tolerance of perovskites^{10,11}.

II. METHOD

Density functional theory^{12,13} (DFT) calculations were performed using the Vienna ab initio simulation package¹⁴ (VASP) and projector augmented-wave potentials^{14–16}. A Perdew-Burke-Ernzerhof¹⁷ (PBE) gradient approximation for the exchange-correlation functional

was chosen in combination with the Grimme *et al.*¹⁸ (D3) correction to capture long-range van der Waals interactions.

The Pb–I framework was kept fixed at experimental positions and lattice parameters¹. This implies that the same lattice constant and atomic distortions apply to all compositions, including the stoichiometric phase of MAPbI₃ [see supporting information (SI), Table SI]. The experiment¹ indicates only a small increase of the lattice parameters for d-MAPbI₃. Cesium atoms were used instead of MA and HEA cations (Fig. 2) to reduce computational efforts. This approach is justified as long as we preserve experimental positions for the lead-iodide framework, which determines the band gap and dispersion near to the band edges. In order to ensure that the conclusions drawn in this paper are not artifacts of the approximation made regarding the structure of d-perovskites, we investigate a structure where HEA⁺ organic cation is substituted in place of (Pb–I)⁺ unit in tetragonal MAPbI₃. Calculations for d-MAPbI₃ with HEA (without Cs and including relaxation of atomic positions for the entire structure) are presented in SI (Figs. S1 and S2). Overall trends are the same (including the type-I band alignment, see Fig. 1), but effects are weaker (smaller band gap opening, smaller band discontinuities, and not so large effective mass enhancement).

When relaxing positions of Cs atoms, the cutoff energy E_{cut} for the plane-wave expansion was increased by 25% above a recommended maximum value in pseudopotential files (VASP tag PREC=High, $E_{\text{cut}} = 275$ eV). In band structure calculations, E_{cut} corresponded to the recommended maximum value in pseudopotential files (VASP tag PREC=Normal, $E_{\text{cut}} = 220$ eV). The absolute change in the band gap as a result of different E_{cut} is marginal (not more than 5 meV). The Monkhorst and Pack¹⁹ shifted k mesh was used for the structural relaxation (Cs atoms only), while the Γ -centred mesh was selected for calculations of the band gap and optical properties (see SI, Table SI). Linear optical properties were calculated within an independent particle approximation ignoring local field effects. The spin-orbit coupling (SOC) was taken into account. The number of bands was increased twice from its default value when calculating optical properties. Other relevant parameters are summarized in Table SI (see SI).

Effective masses were determined at the PBE+SOC level by fitting the energy dispersion $E(k)$ to a second order polynomial within the range of $\Delta k = \pm 0.003 \text{ \AA}^{-1}$ near to the band edge sampled with 9 – 11 intermediate k points.

Band structure unfolding and Bloch character calculations were done with the fold2Bloch

code²⁰ following the method outlined in Ref. 21. The code computes spectral weights $w_n(\mathbf{k})$ in the range 0–1 for each n 's energy eigenstate to reflect its Bloch \mathbf{k} character. The number of spectral weights computed per eigenstate corresponds to the multiplicity of the supercell. The approach to the calculation of spectral weights is based on remapping the supercell reciprocal space with a mesh that is compatible with the translational symmetry of a primitive cell. Computed Bloch spectral weights for each eigenstate in a supercell calculation fulfill the normalization $\sum_{\mathbf{k}} w_n(\mathbf{k}) = 1$, which implies that the cumulative probability of finding an electron in any of multiple \mathbf{k} 's states adds up to one. Further details of the method and implementation can be found elsewhere^{20,21}.

Vesta²² was used for visualization of atomic structures. Structure files used in this work can be accessed at the Cambridge Crystallographic Data Center (CCDC) under deposition No. 1962341-1962347 and 1962349.

III. RESULTS AND DISCUSSION

A. Structure of d-perovskites and defect formation enthalpy

It is worth commenting on the structure of d-perovskites before going into discussion of their electronic properties. Here we adopt the structure proposed in Ref. 1, where the lattice parameters and coordinates of Pb–I framework are specified. Fig. 2 shows the stoichiometric structure as well as two d-perovskite structures that differ by the I/Pb ratio (3.22 and 3.5) corresponding to the defect concentration of $x = 0.1$ and 0.2 , respectively. The closer the ratio to its stoichiometric value of 3, the less deficient is the perovskite structure. The ratio of I/Pb = 3.5 is at the upper limit of deficiency attained experimentally¹, if not hollow. In the sections that follows, Cs atoms are used instead of organic cations (MA or HEA) for computational performance reasons (see also the case study with HEA in Method section).

The self-compensation defect formed by A cation substituted for $M-X$ unit might also be unintentionally present in smaller concentrations in stoichiometric AMX_3 perovskite structures. To evaluate a formation energy of the corresponding defect [Eq. (1)] we built a $2 \times 2 \times 2$ supercell of d-MAPbI₃ with and without defect as shown in Fig. S1 (see SI). The defect structure corresponds to the larger HEA⁺ cation substituted in place of a Pb–I unit

of MAPbI_3 . The corresponding reaction can be expressed as



Giving the size of the supercell, it corresponds to a defect concentration of $x = 0.125$. We assign the DFT total energy difference between reactants and products to the defect formation enthalpy

$$\Delta H_d = E_{\text{MA}_8\text{Pb}_8\text{I}_{24}} + E_{\text{PbI}_2} - (E_{\text{MA}_8(\text{HEA})\text{Pb}_7\text{I}_{23}} - E_{\text{HEAI}}). \quad (3)$$

The calculations yield $\Delta H_d = 0.78$ eV. This result can be compared with another common Schottky defect in perovskites, namely PbI_2 vacancy. The later has a formation energy of about 0.2 eV²³ suggesting that the defect described by Eq. (1) is less likely to form spontaneously.

B. Band gap opening, band alignment, dual band gap and optical properties

There is a significant opening of the band gap in the deficient structures relative to the reference (stoichiometric) perovskite. Computed band gaps are listed in Table I. The band gap opening is more prominent when relativistic effects are included. The calculated band gap increment (0.74 eV with SOC and 0.36 eV without SOC) is within the range of the experimental band gap opening quoted in Sec. I.

It is quite unusual that defects lead to opening of the band gap. It implies that the defect described in Eq. (1) does not create states within the band gap. The Pb–I vacancy introduces five undercoordinated I atoms and one Pb atom. A simplified orbital energy diagram of halide perovskites in Fig. 3(a) is intended to rationalize this result. Since both top of the valence band and bottom of the conduction band are formed by antibonding orbitals^{24,25}, compensated dangling bonds of Γ^- are expected to appear *within* the bulk of valence band states as confirmed by the analysis of atom-specific contribution to energy eigenstates (see SI, Fig. S3).

Opening of the band gap can be viewed as a confinement effect. Defects form local regions of $A'X$ with a wide band gap and confine electronic states within the perovskite region as shown in Fig. 3(b,c). The higher the concentration of defects, the smaller the distance between $A'X$ regions leading to a higher confinement energy. Since it is a confinement

effect, we would expect it to be operational at a high enough concentration of defects when the distance between them becomes comparable to the lattice spacing.

It is interesting to see whether this type of defects will act as trap for charge carriers. The greater band gap of d-perovskites alone does not provide an answer to this question, as there are different types of band alignment possible (namely type I or II, Fig. 1). The band alignment calculations were performed using core Cs-5s states as an energy reference to relate position of the valence and conduction bands in different calculations. (See Refs. 26,27 for details of this method as well as discussion of its accuracy and validation.) Our calculations suggest an energy misalignment of approximately $\delta_v = 0.1$ eV (Table I) between the VBE of stoichiometric and d-perovskite with VBE of d-perovskite being lower in energy as shown in Fig. 4(a). This result is consistent with the recent data on d-FAPbI₃⁶. Most of the band gap difference between stoichiometric and d-perovskite comes from the CBE offset δ_c (Table I). This result indicates a type-I band alignment between stoichiometric and d-perovskites. Thus, Pb-I deficient regions should pose no threat for trapping of charge carriers. It is an interesting result, since defects are generally expected to create either shallow or deep traps within the band gap, which is not the case here.

The disparity between the CBE and VBE band offset energies (δ_c and δ_v) can be attributed to a much greater confinement energy for electrons than holes in the $A'X/AMX_3$ heterostructure shown in Fig. 3(b,c). The diagram was constructed based on experimental data for an electron affinity of 0.63 eV²⁸ and an ionization energy of 7.1 eV²⁹ for CsI as well as an electron affinity of 4.4 eV and an ionization energy of 5.9 eV for MAPbI₃²⁵.

The band diagram in Fig. 4(a) can now be used to explain an apparent dual band gap in reflectance measurements [see Fig. 2(a) in Ref. 1 also shown schematically in Fig. 4(b)]. The dual band gap character becomes more prominent with increasing deficiency of the perovskite. The first and second experimental gaps amount to $E_g^{(1)} \approx 1.6$ eV and $E_g^{(2)} \approx 2.0 - 2.1$ eV, respectively.

To further inspect the possible origin of the dual band gap, we compute the optical response of d-perovskites. Calculations of optical properties (Fig. 5) indicate an anisotropy of the dielectric function in d-perovskites. The anisotropy becomes more prominent as the deficiency increases. The onset of absorption in d-perovskites is shifted to higher energies following the band gap trend and in accord with experimental observations^{1,3}. The calculated imaginary part of the dielectric function (Fig. 5) shows a clear blue shift, but no features that

can be interpreted as the second gap. The monotonic behavior could be attributed to the growing density of states as the energy moves away from band edges into the bands (see SI, Fig. S4). Thus, to rationalize experimental observation, one needs to assume a coexistence of stoichiometric MAPbI₃ regions that contribute to $E_g^{(1)}$ and Pb–I deficient regions with higher optical transition energies at $E_g^{(2)}$ as shown in Fig. 4(a).

C. Carrier transport properties

Since d-perovskites are considered as a potential absorber material in solar cells,^{1,5} it is interesting to explore how deficiencies affect intrinsic transport characteristics of this material. The band dispersion near CBE and VBE is shown in Fig. 6. As a general trend, d-perovskites show less dispersive bands, which was also noticed previously³. The band dispersion is mostly affected within the x – y plane. It is the plane perpendicular to direction of the tunnel formed after removing the Pb–I unit [Fig. 2(c)]. At high concentration of defects, the dispersion remains only along z axis that electronically resembles 1D structures.

Effective mass calculations (Table II) corroborate conclusions drawn from the band dispersion data in Fig. 6. The effective mass magnitude $|m^*/m_e|$ in the stoichiometric perovskite is of the order of 0.09 for electrons and holes. Our calculated effective masses are about twice less than previous calculations for CsPbI₃, $|m^*/m_e| = 0.15 - 0.20$ ^{30,31} and for MAPbI₃ $|m^*/m_e| = 0.13 - 0.22$ ³², and are also underestimated when compared to more accurate results based on many body perturbation theory within the GW approximation^{33,34}. It is expected that the bare DFT-PBE will underestimate effective masses³⁵ due to the band gap error. Meanwhile, our computed reduced effective masses (0.08 – 0.09) compare quite well to the experimental exciton reduced mass of MAPbI₃ $0.104|m_e|$ ³⁶. In d-perovskites, the effective masses of electrons and holes increase as the material becomes more defective. At the intermediate concentration of defects (the I/Pb ratio of 3.22, $x = 0.1$), the increase is moderate (by a factor of 4 at most). The effective mass becomes anisotropic. The lightest effective mass corresponds to the band dispersion along channels. The charge transport (especially holes) becomes effectively 1D when the concentration of defects (Pb–I vacancies) approached the upper limit (I/Pb = 3.5, $x = 0.2$) and defects form a hollow tunnel along z axis.

D. Band unfolding of d-perovskite

To gain further insight into changes of the band structure introduced by Pb–I deficiency we employ a band unfolding technique. It requires the supercell to have lattice vectors aligned with lattice vectors of the primitive cell. Here we use the quasi-cubic cell (blue area on Fig. 7) as a primitive unit for projecting an unfolded band structure. We created a supercell in such a way that its lattice vectors are aligned with those of the primitive cubic cell while preserving the periodicity of the deficient perovskite structure at the same time (Fig. 7). The smallest supercell that fulfills these conditions is shown as a pink area in Fig. 7. Its multiplicity factor is $5 \times 5 \times 2$ relative to the cubic primitive cell.

The unfolded band structures are shown in Fig. 8. A band structure of the stoichiometric phase was also unfolded for comparison. It shows a direct band gap at R point [Fig. 8(a)] as expected from previous studies of 3D lead-halide perovskites^{24,37}. The Bloch character of electronic states near to the band edges of the stoichiometric perovskite is almost 100%. In the d-perovskite with high density of defects (I/Pb = 3.5, $x = 0.2$), 70 – 80% of the Bloch character is preserved [Fig. 8(b)]. The band structure remains direct at R point. The band dispersion characteristics reflect the above discussion of effective masses. A well-preserved Bloch character near R point implies that states are not localized despite the presence of structural defects. This is further corroborated when using inverse participation ratio³⁸ of electronic eigenstates as a measure of localization (see SI for details). In fact, analysis of the inverse participation ratio (see SI, Fig. S5) did not reveal any localized states near to the band edges either. This unusual behavior of d-perovskites is another manifestation of a defect tolerance inherent to hybrid halide perovskites^{10,39}.

IV. CONCLUSION

The electronic structure evolution of d-perovskites with a general formula $(A, A')_{1+x}M_{1-x}X_{3-x}$ was investigated using the density functional theory for an intermediate and a high concentration of defects ($x = 0.1$ and 0.2 , respectively). The formation enthalpy of Pb–I deficient defects filled with a larger HEA⁺ cation for charge compensation is about 0.8 eV, which is relatively large when compared to other native defects in halide perovskites. Our calculations reproduce the opening of a band gap also observed in previous theoretical

studies. The band gap opening is achieved mostly due to moving the conduction band edge upward in energy. In addition, the valence band edge shifts slightly toward lower energies by approximately 0.1 eV. This leads to the formation of a type-I band alignment between d-perovskites and the stoichiometric phase. As a result, Pb–I deficient regions do not act as traps for charge carriers. This result is rationalized via a confinement model where the defect is viewed locally as a wide band gap region of $A'X$.

Presence of this type of defects in moderate quantities ($x \leq 0.1$) should not hinder the ability of perovskite material to function as a good solar cell absorber. Based on effective mass calculations, we expect a reduced mobility of electrons and holes in d-MAPbI₃ ($x = 0.1$) by a factor of 4 at most. However, at high concentration of defects ($x = 0.2$), the charge transport ceases within a lateral plane due to an extremely heavy effective mass of holes thereby posing an issue for extracting photogenerated charge carriers. Interestingly, no localized states were detected near to the band edges of d-perovskites. It is attributed to the fact that both conduction and valence band edges are derived from antibonding orbitals.

Investigation of the optical properties of d-perovskites reveal no anomalies in absorption spectra. The absorption spectrum increases monotonously up to a photon energy of about 2.5 eV. Thus, dual band gap features observed experimentally at a high concentration of defects ($x > 0.15$) are attributed to coexistence of deficient and stoichiometric regions.

Finally, we would like to comment on the electronic structure dimensionality of d-perovskites. Deficient perovskites with a moderate density of defects ($I/Pb = 3.22$, $x = 0.1$) behave as a 3D material with anisotropic effective masses. As the density of defects increases ($I/Pb = 3.5$, $x = 0.2$), hollow regions connect into tunnels along the z axis, which remains the only direction for efficient band dispersion leading to a 1D electronic structure. Also, the band gap opening of d-perovskites exceeds that of 2D perovskites. This statement builds on the fact that a PL emission wavelength of 2D 1 ML MAPbI₃ perovskite is centered at 720 nm⁴⁰, while d-MAPbI₃ have the secondary gap at shorter wavelengths (about 600 nm¹).

SUPPLEMENTARY MATERIAL

See the supplementary material for parameters used in electronic structure calculations; structure and band dispersion results for d-MAPbI₃; density of states, partial density of states, and the inverse participation ratio analysis. References cited therein are^{38,41,42}.

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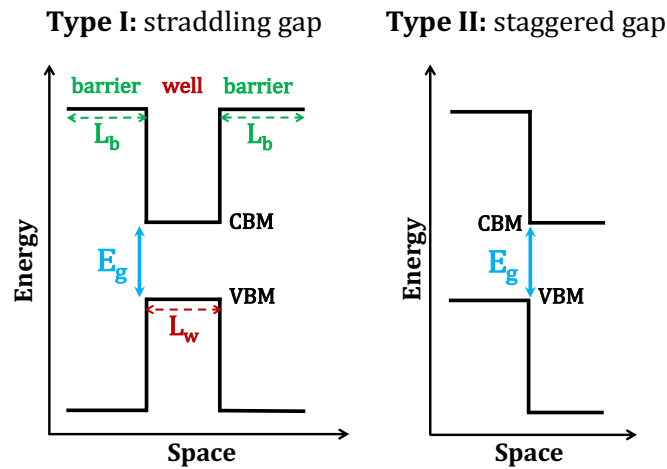


FIG. 1: Band offsets at a semiconductor heterojunction illustrated for two types of band alignment. In type I, the conduction band of the well is lower than that of the barrier. Oppositely, the valence band of the well lies higher than that of the barrier. In type II, both the conduction and valence bands of one of the two materials building the heterojunction are located below those of the second material.

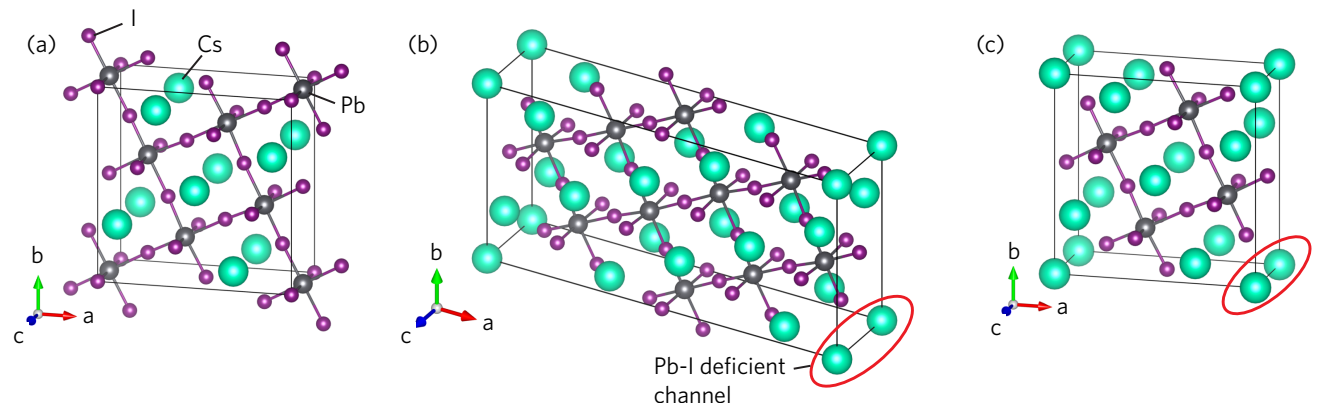


FIG. 2: (a) Stoichiometric perovskite structure without deficiencies, (b) deficient perovskite with I/Pb ratio of 3.22 ($x = 0.1$), (c) I/Pb ratio of 3.5 ($x = 0.2$). Cesium atoms are used instead of MA cations. Defects are arranged in continuous channels along c axis.

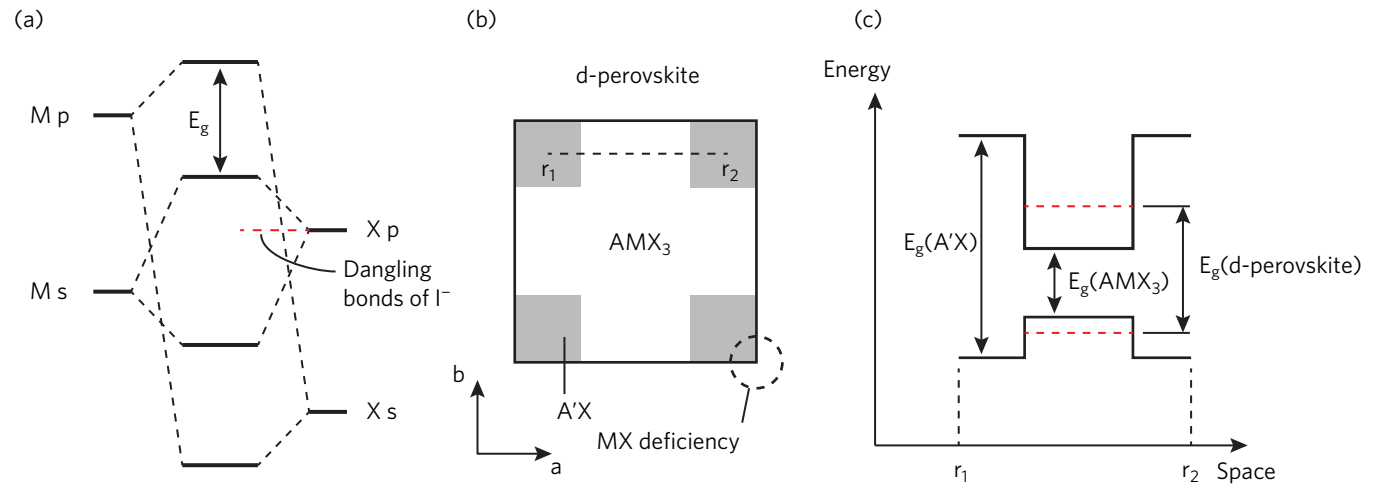


FIG. 3: (a) Simplified orbital energy diagram of AMX_3 . Only states of the MX_3^- lattice are shown. SOC effects as well as $s-s$ and $p-p$ hybridizations are omitted. (b) Schematic structure of d-perovskite. Gray regions correspond to the local $A'X$ chemical environment. (c) Band diagram plotted along the $r_1 - r_2$ path. Electronic states of d-perovskite are derived from AMX_3 confined between regions of $A'X$. Since holes are less confined than electrons, the corresponding confinement energies are ordered as $\delta_c > \delta_v$.

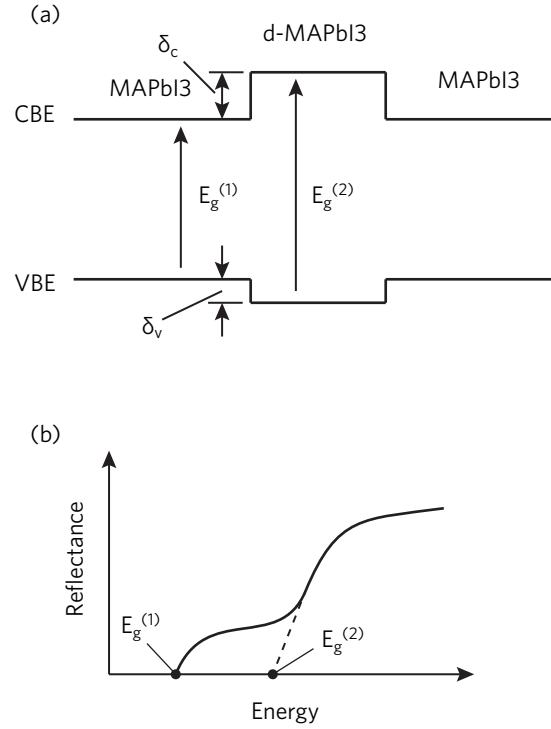


FIG. 4: (a) Band alignment diagram that shows a d-perovskite region embedded in the stoichiometric MAPbI₃. (b) Photoreflectance spectrum (schematic) that shows two apparent band gaps, that originate from two optical transitions $E_g^{(1)}$ and $E_g^{(2)}$ shown on panel (a).

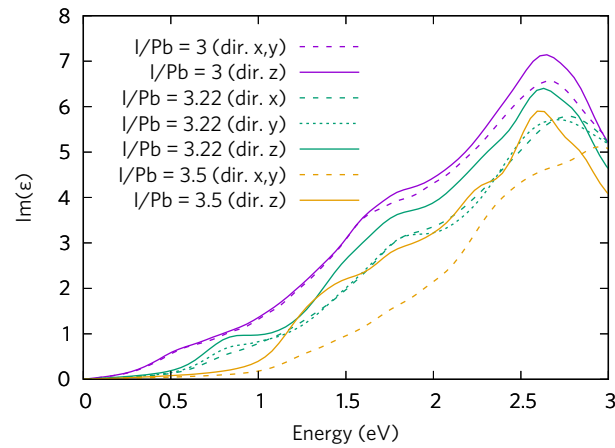


FIG. 5: Frequency-dependent imaginary part of the dielectric function for a stoichiometric perovskite structure without deficiencies ($x = 0$, violet) and d-perovskites with I/Pb ratio of 3.22 and 3.5 ($x = 0.1$, green, and 0.2, orange, respectively). Only diagonal components of the dielectric tensor are presented (ϵ_{xx} , ϵ_{yy} and ϵ_{zz}); off-diagonals components are close to zero. The Cartesian directions correspond to the coordinate system in Fig. 2.

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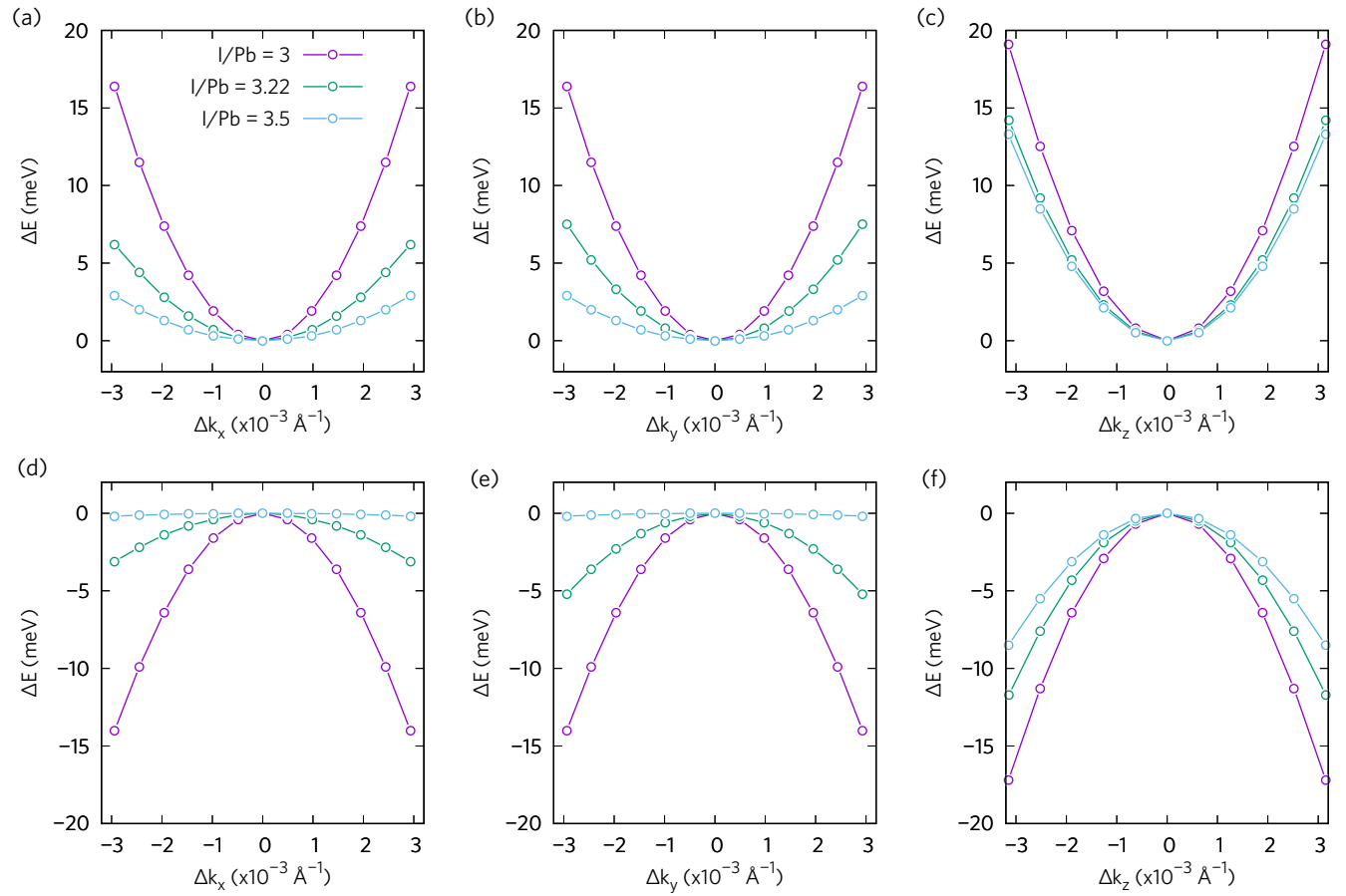


FIG. 6: Band dispersion along x , y , and z direction in reciprocal space in the vicinity of (a-c) CBE and (d-f) VBE for a stoichiometric perovskite structure without deficiencies ($I/Pb = 3$, $x = 0$) and deficient perovskites with I/Pb ratios of 3.22 and 3.5 ($x = 0.1$ and 0.2, respectively). Bands become less dispersive with increasing the deficiency x in particular in the lateral directions (x and y).

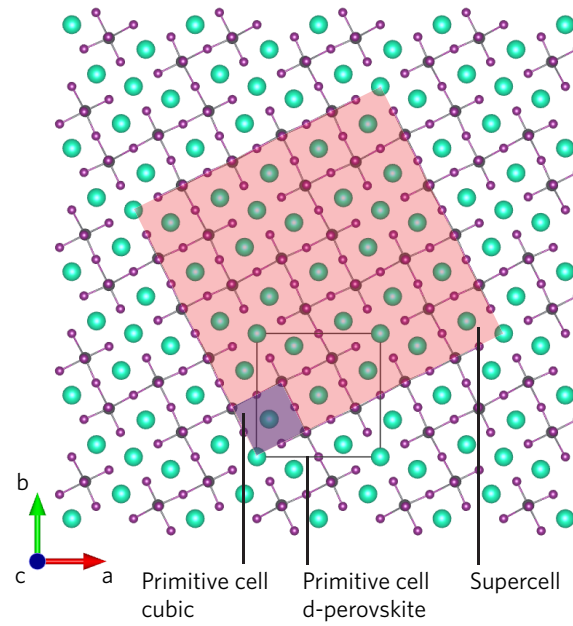


FIG. 7: Relation between the d-perovskite unit cell (black solid line), the cubic primitive cell (blue area), and the supercell (pink area).

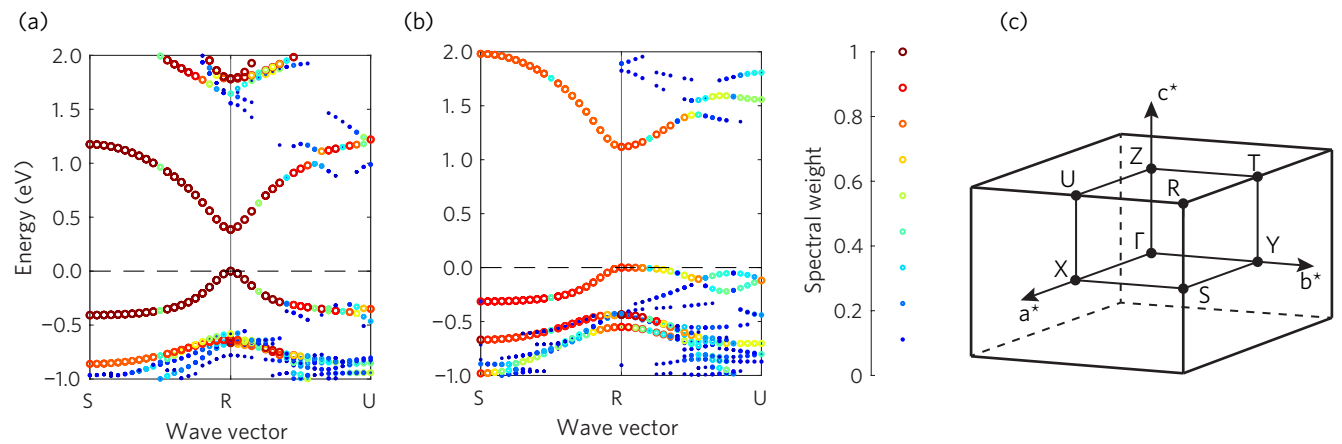


FIG. 8: Effective band structure of (a) stoichiometric and (b) d-perovskite ($I/Pb = 3.5$, $x = 0.2$). The color and size of symbols correspond to the Bloch spectral weight. Points with the spectral weight less than 0.05 are not shown. (c) Brillouin zone of an orthorhombic lattice. The Bloch character is well preserved (greater than 70%) in d-perovskites despite of less dispersive bands.

TABLE I: Band gaps and band alignments (eV)

Parameter	Structure		
	stoichiometric Fig. 2(a)	I/Pb = 3.22, $x = 0.1$ Fig. 2(b)	I/Pb = 3.5, $x = 0.2$ Fig. 2(c)
E_g PBE (no SOC)	1.47	1.59	1.83
E_g PBE (with SOC)	0.38	0.66	1.12
δ_v PBE (with SOC) ^a	n/a	0.13	0.13
δ_c PBE (with SOC) ^a	n/a	0.15	0.61

^a See Fig. 4(a) for definition of band alignments. The positive value of δ_v indicates that the VBE is shifted down in energy relative to the stoichiometric phase. The positive value of δ_c corresponds to the upward shift of the CBE in energy.

TABLE II: Effective masses of electron and holes m^*/m_e .

Carrier type, direction	Structure		
	stoichiometric Fig. 2(a)	I/Pb = 3.22, $x = 0.1$ Fig. 2(b)	I/Pb = 3.5, $x = 0.2$ Fig. 2(c)
Electrons			
$R \rightarrow S$	0.08	0.11	0.11
$R \rightarrow U$	0.08	0.17	0.45
$R \rightarrow T$	0.08	0.21	0.45
Holes			
$R \rightarrow S$	-0.09	-0.13	-0.18
$R \rightarrow U$	-0.09	-0.25	-7.4
$R \rightarrow T$	-0.09	-0.42	-7.4

Supporting information:

Electronic properties of Pb-I deficient lead halide perovskites

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TABLE SI: Parameters used in electronic structure calculations.

Parameter	Structure	
	Stoichiometric perovskite [I/Pb = 3, Fig. 2(a)] and d-perovskite [I/Pb = 3.5, Fig. 2(c,d)]	d-perovskite [I/Pb = 3.22, Fig. 2(b)]
Cs atomic positions relaxation:		
$a = b$ (Å)		14.3120
c (Å)	6.3656	12.7312
valence electrons of Cs		$5s^2 5p^6 6s^1$
valence electrons of Pb		$6s^2 6p^2$
valence electrons of I		$5s^2 5p^5$
positions of Pb and I		fixed
positions of Cs		relaxed (max force 0.02 eV/Å)
E_{cut} (eV)		275
k mesh	$2 \times 2 \times 4$ (shifted Monkhorst-Pack)	$2 \times 2 \times 2$ (shifted Monkhorst-Pack)
SOC		disabled
Energy convergence (eV)		10^{-7}
Band gap:		
E_{cut} (eV)		220
k mesh	$2 \times 2 \times 4$ (Γ -centred)	$2 \times 2 \times 2$ (Γ -centred)
SOC		enabled
Optics:		
E_{cut} (eV)		220
k mesh	$12 \times 12 \times 24$ (Γ -centred)	$12 \times 12 \times 12$ (Γ -centred)
SOC		enabled
number of bands	400	800

In order to ensure that the conclusions drawn in this paper were not artifacts of the approximation made regarding the structure of d-perovskites, we investigate a defect where HEA^+ organic cation is substituted in place of $(\text{Pb-I})^+$ unit in tetragonal MAPbI_3 as shown in Fig. S1. This structure corresponds to the ratio $\text{I/Pb} = 3.3$ ($x = 0.125$), which corresponds to a moderate defect density. A 300 K stoichiometric tetragonal MAPbI_3 structure was taken from Ref. 1. Atomic positions were relaxed in both structures (stoichiometric and d- MAPbI_3). The nitrogen 2s states were used as a reference for band alignment calculations. Observations include:

- band gap opening by 0.23 eV (calculated at PBE+SOC level, similar to results in Table I);
- the band gap of tetragonal d- MAPI_3 structure remains direct at Γ point;
- the stoichiometric and d- MAPbI_3 form a type-I band alignment with the valence band discontinuity of $\delta_v = 0.06$ eV (twice less than in Table I).

The band structure of d- MAPbI_3 (Fig. S2) shows a slightly less dispersive character as compared to its stoichiometric counterpart. However, the increase in the effective mass are not drastic (only about 50%), which is less than data for unrelaxed structures with Cs (Fig. 6 and Table II) for the equivalent I/Pb ratio.

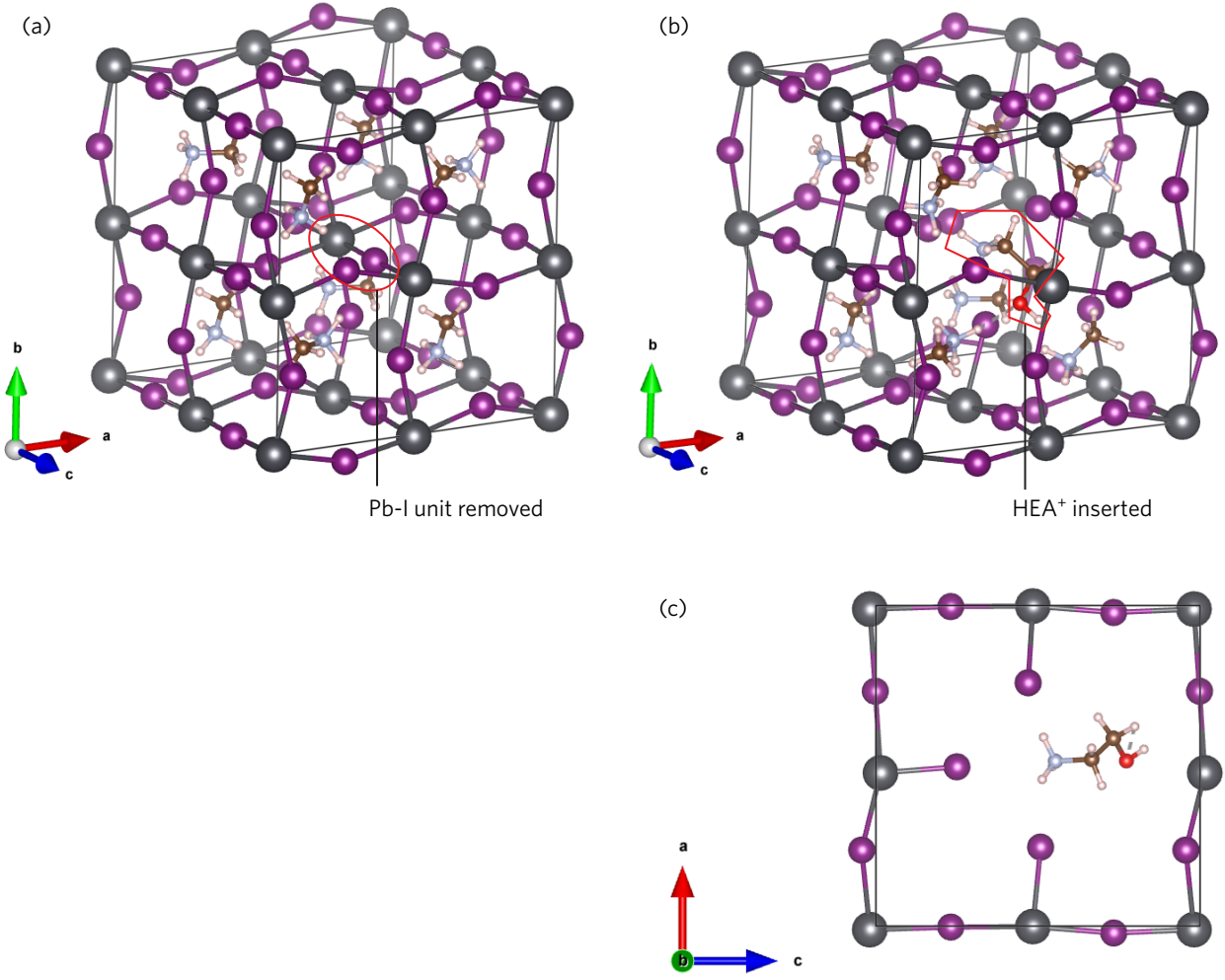


FIG. S1: (a) Stoichiometric tetragonal MAPbI₃. (b) d-MAPbI₃ with the ratio I/Pb = 3.3 ($x = 0.125$). (c) Arrangement of HEA⁺ in Pb-I framework. HEA⁺ fits well in the cavity forming hydrogen bonds with neighbor MA⁺ cations and causes only a small stress of a few kBar.

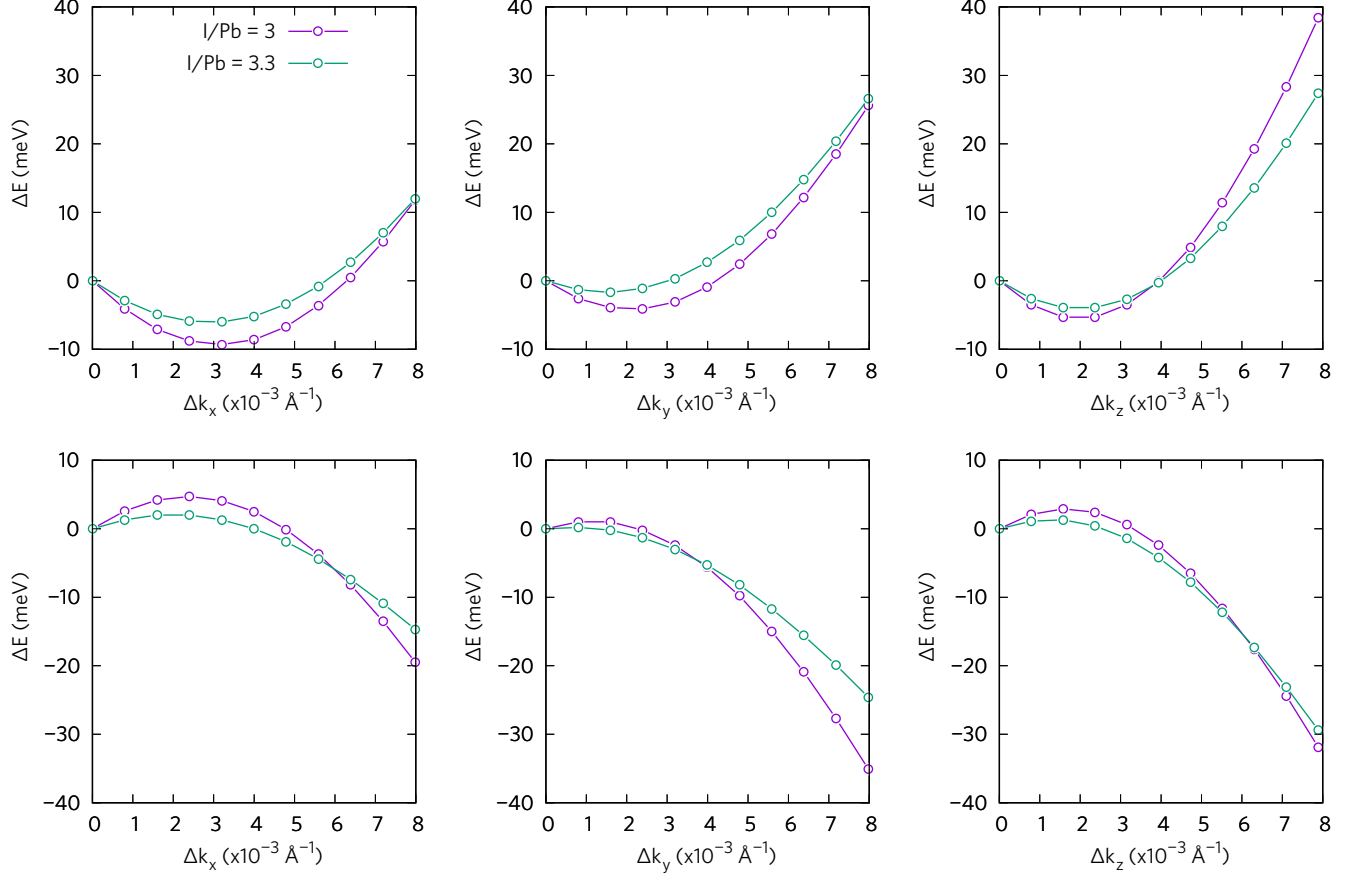


FIG. S2: Band dispersion along x , y , and z direction in reciprocal space in the vicinity of (a-c) CBE and (d-f) VBE for a stoichiometric MAPbI_3 without deficiencies ($I/\text{Pb} = 3$) and d- MAPbI_3 with HEA [$I/\text{Pb} = 3.3$, $x = 0.125$, Fig. S1(b)]. Bands of d-perovskite become less dispersive, but the change is not as dramatic as in the model structure with Cs (Fig. 6).

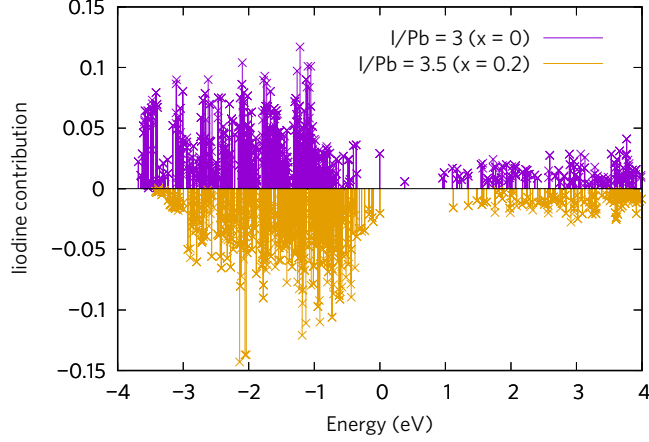


FIG. S3: Contribution of undercoordinated iodine atoms to energy eigenstates within ± 4 eV from the band edges. Negative values are used to the vertical axes to facilitate the comparison. Origin of the energy scale is set at the top of the valence band. The probability of finding an electron in the vicinity of undercoordinated iodine atoms is shifted to higher energies in the valence band, while values in the proximity of the band edges remain unchanged, which indicates no dangling bonds.

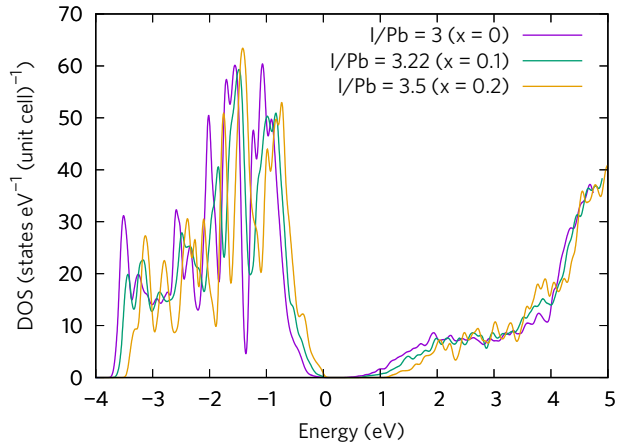


FIG. S4: Density of states for a stoichiometric perovskite structure without deficiencies ($x = 0$) and d-perovskite with I/Pb ratio of 3.22 and 3.5 ($x = 0.1$ and 0.2 , respectively). Origin of the energy scale is set at the top of the valence band.

To complement data on unfolding of the band structure of d-perovskites we performed additional search for localized states formed near to the band edges. Here we employ an inverse participation ratio (IPR) as a measure of localization. The IPR for each eigenvalue E_i has been calculated based on the probabilities p_n of finding an electron within a muffin-tin sphere of an atom n^2

$$\text{IPR}(E_i) = \frac{\sum_n p_n^2(E_i)}{[\sum_n p_n(E_i)]^2}. \quad (\text{S1})$$

Here the summation index n runs over all atoms in the simulation cell. The participation ratio represents a number of atomic sites which confine the wave function $\psi_i(\mathbf{r})$. Thus, the lower limit of the IPR is the inverse number of atoms in the cell that corresponds to a completely delocalized state. The upper upper bound is $\text{IPR} = 1$, which corresponds to a state fully localized at a single atom. Data for $p_n(E_i)$ can be found in PROCAR file generated by enabling the VASP tag LORBIT=11. The PROCAR file was processed with "prPROCAR.m" Octave code, which is a part of a "VASPtools" subroutine collection³. Spectrally-resolved IPR results are presented in Fig. S5. They show no notable signatures of localization near to the band edges when comparing d-perovskites with the stoichiometric structure.

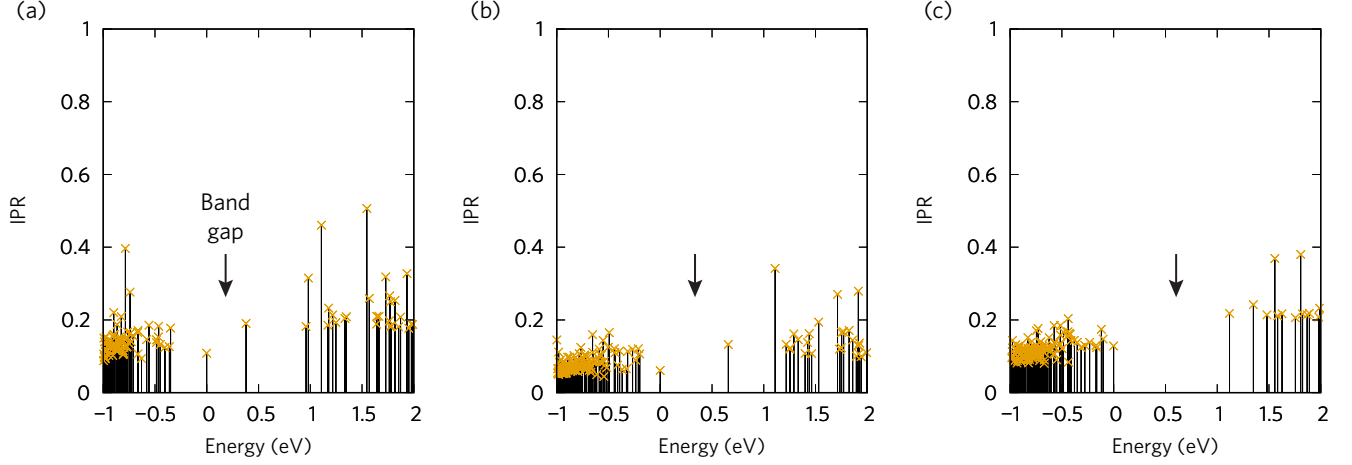


FIG. S5: Inverse participation ratio (IPR) vs energy of electronic states in (a) stoichiometric, (b) d-perovskite ($I/Pb = 3.22$, $x = 0.1$), and (c) d-perovskite ($I/Pb = 3.5$, $x = 0.2$). The lower limit of $IPR = 0.05$, 0.026 , and 0.056 correspond to fully delocalized states in Pb–I framework (different values are due to different number N of Pb and I atoms in each structure, $\min(IPR) = N^{-1}$). The upper limit of $IPR = 1$ corresponds to the extreme localization of the wave function in the vicinity of a single atom. Origin of the energy scale is set at the top of the valence band. No localized states are observed at the band edges even in d-perovskite with the high defect concentration x .

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